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MODELLING OF SUBMERGED  
CABLE DYNAMICS

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# ABSTRACT

Results from a series of simulated submerged cable maneuvers are presented. The simulations are obtained using a three-dimensional, finite-segment model of the cable. The model, called UCIN-CABLE, consists of a series of ball-and-socket connected rigid rods. Fluid drag, inertia and buoyancy forces are included.

Two types of simulation are presented: buoy release and anchor drop. The results compare favorably with experimental data and with data obtained from finite-element modelling.

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# NOTATION

$a$	Cross section area of $L_j$ .
$a, b, c$	Coefficients defined by Equations (26), (27) and (28).
$a_{GN}$	Component of fluid acceleration normal to $L_j$ at $G_j$ . [See Equation (24).]
$a_N$	Normal component of fluid acceleration relative to $L_j$ at P.
$a_{W/G}$	Acceleration of the fluid relative to $G_j$ .
$a_{lp}$	Mass matrix. [See Equations (1) and (2).] <span style="float: right;">(<math>l, p=1, \dots, n</math>)</span>
$A, B, C$	Coefficients defined by Equations (5), (6) and (7).
$b_N$	Buoyancy force on $L_j$ at P, per unit length.
$B_N$	Equivalent buoyancy force on $L_j$ at $G_j$ . [See Equation (19).]
$C_M, C_N, C_T$	Coefficients defined by Equations (8), (9), and (10).
$d$	Diameter of $L_j$ .
$e_{mnk}$	Permutation symbol.
$f$	Applied force on $L_j$ at P. [See Equation (4).]
$f_l$	Defined by Equation (3). <span style="float: right;">(<math>l=1, \dots, n</math>)</span>
$F_j$	Equivalent applied force on $L_j$ at $G_j$ . [See Equation (15).]
$F_{jk}$	$n_k$ components of $F_j$ . <span style="float: right;">(<math>j=1, \dots, N; k=1, 2, 3</math>)</span>
$F_M$	Equivalent added mass force on $L_j$ at $G_j$ . [See Equation (17).]
$F_N$	Equivalent normal drag force on $L_j$ at $G_j$ . [See Equation (18).]
$F_T$	Equivalent tangential drag force on $L_j$ at $G_j$ . [See Equation (19).]
$F_l$	Generalized active forces. <span style="float: right;">(<math>l=1, \dots, n</math>)</span>
$g$	Gravity constant.
$G_j$	Mass center of $L_j$ . <span style="float: right;">(<math>j=1, \dots, N</math>)</span>
$I_{jkh}$	$n_k$ components of the inertia dyadic of $L_j$ with respect to $G_j$ . <span style="float: right;">(<math>j=1, \dots, N; k, h=1, 2, 3</math>)</span>
$k$	Vertical (up) unit vector.
$L$	Length of $L_j$ .
$L_j$	Typical cable segment. <span style="float: right;">(<math>j=1, \dots, N</math>)</span>
$m_j$	Mass of segment $L_j$ .
$n$	Number of degrees of freedom.
$\underline{n}$	Unit vector parallel to $L_j$ .
$\underline{n}_k$	Unit vectors fixed in R. <span style="float: right;">(<math>k=1, 2, 3</math>)</span>
$N$	Number of cable segments.
$P$	A typical point on $L_j$ .

$R$	An inertial reference frame.
$R_{eN}, R_{eT}$	Reynolds numbers defined by Equations (11) and (12).
$T_j$	Equivalent applied torque on $L_j$ . [See Equation (16).]
$T_{jk}$	$n_k$ components of $T_j$ . ( $j=1, \dots, N; k=1, 2, 3$ )
$T_M$	Equivalent added mass torque on $L_j$ . [See Equation (22).]
$T_N$	Equivalent normal drag torque on $L_j$ . [See Equation (23).]
$v_{jpk}$	$n_k$ components of the partial velocity of $G_j$ . ( $j=1, \dots, N; p=1, \dots, n; k=1, 2, 3$ )
$v_{GN}$	Components of fluid velocity normal to $L_j$ at $G_j$ . [See Equation (25).]
$v_N$	Normal component of fluid velocity relative to $L_j$ at P.
$v_T$	Tangential component of fluid velocity relative to $L_j$ at P.
$v_{W/G}$	Velocity of the fluid relative to $G_j$ .
$\bar{w}$	Weight of $L_j$ per unit length.
$\bar{w}$	Weight of $L_j$ . [See Equation (20).]
$X_p$	Orientation angles. ( $p=1, \dots, n$ )
$\mu$	Viscosity of the fluid.
$\rho$	Fluid mass density.
$\rho_c$	Cable mass per unit length.
$\omega_{jpk}$	$n_k$ components of the partial angular velocity of $L_j$ . ( $j=1, \dots, N; p=1, \dots, n; k=1, 2, 3$ )

## INTRODUCTION

In a series of recent papers [1,2,3]\*, we presented a method for modelling cable dynamics. The method is based upon previously developed general procedures for finite segment modelling of multibody systems [4,5]. In this report we present results obtained by applying the method in a series of submerged cable configurations. The results are compared with experimental data reported by Palo at the Navy's Civil Engineering Laboratory [6], and with numerical data obtained from SEADYN [7]--a finite element cable computer code.

A brief review of the method itself is presented in the first section of the report.

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\*Numbers in brackets refer to References at the end of the report.

## I. FINITE SEGMENT MODELLING

We model a cable by a series of ball-and-socket connected rigid links, or segments, as depicted in Figure 1. The dimensions and physical parameters of the segments are arbitrary. Also, we let the segments be subjected to general force fields. This allows us to simulate an arbitrary fluid environment, as well as gravity and buoyancy forces. Finally, cable stiffness or elasticity, if it is significant, can be modelled by springs and dampers between the segments.

We describe the orientation and configuration of the system using the relative angles between the segments. We assume that in a specific maneuver that one end of the cable may be attached to a towed body and that the motion of the other end is specified (for example, it may be fixed). We assume that the initial configuration of the system is known. The objective of the analysis is then to define the subsequent configuration of the system.

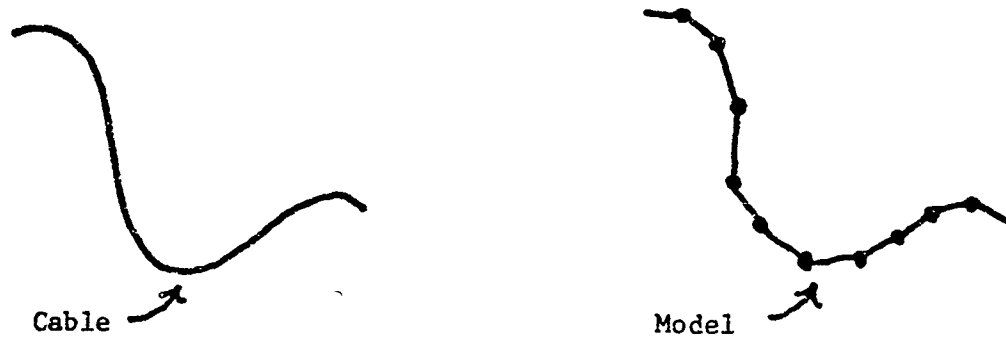


Figure 1. A Finite Segment Model of a Cable.

Using procedures developed for finite segment modelling of multi-body systems [4,5], we can obtain explicit governing dynamical equations of motion for the cable model. What is more, we can obtain these equations in a form suitable for conversion into algorithms for numerical integration. Specifically, the equations may be written in the form:

$$a_{lp} \ddot{X}_p = f_l \quad (l, p = 1, \dots, n) \quad (1)$$

where the  $X_p$  ( $p=1, \dots, n$ ) represent the relative orientation angles of the cable segments, where  $n$  is the number of degrees of freedom. (Regarding notation, there is a sum for repeated indices over the range of the index.)  $a_{lp}$  and  $f_l$  are given by the expressions:

$$a_{lp} = m_j v_{jpk} v_{jlk} + I_{jkh} \omega_{jph} \omega_{jlk} \quad (2)$$

and

$$f_l = F_l - (m_j v_{jlk} v_{jqk} \dot{X}_q + I_{jkh} \omega_{jlk} \dot{\omega}_{jqk} \dot{X}_q + e_{mnk} \omega_{jqn} \omega_{jsr} \omega_{jlk} I_{jmr} \dot{X}_q \dot{X}_s) \quad (3)$$

where  $m_j$  ( $j=1, \dots, N$ ) is the mass of the  $j$ th cable segment  $L_j$ ,  $N$  is the number of cable segments;  $I_{jkh}$  ( $k, h=1, 2, 3$ ) are the components of  $I_j$ , the inertia dyadic of  $L_j$  with respect to its mass center  $G_j$ , referred to unit vectors  $\underline{n}_k$  fixed in an inertial reference frame  $R$ ;  $\omega_{jpk}$  and  $v_{jpk}$  ( $j=1, \dots, N$ ;  $p=1, \dots, n$ ;  $k=1, 2, 3$ ) are the  $\underline{n}_k$  components of the partial angular velocities and partial velocities of  $L_j$  and  $G_j$  [8];  $F_l$  ( $l=1, \dots, n$ ) are the generalized active forces [8] developed from the forces applied to  $L_j$ ; and  $e_{mnk}$  is the permutation symbol [9].

Equations (1) form a set of  $n$  simultaneous nonlinear ordinary differential equations determining the orientation angles  $X_p$ . Since the coefficients  $a_{lp}$  and  $f_l$  are functions of the arrays  $v_{jlk}$  and  $\omega_{jlk}$  and their derivatives, the system of governing equations can automatically be generated once the  $v_{jlk}$  and  $\omega_{jlk}$  arrays and their derivatives are determined. Simple algorithms for computing  $v_{jlk}$  and  $\omega_{jlk}$  and their derivatives have been developed. They are recorded in References [4] and [5].



## II. FLUID AND GRAVITY FORCE MODELLING

The applied forces contributing to the generalized active forces  $F_k$  consist of fluid and gravity forces. The fluid forces may be represented as: 1) Normal drag forces: 2) Tangential drag forces: 3) Added mass forces; and 4) Buoyancy forces. Reference [2] contains a detailed analysis of each of these forces. Specifically, it is shown in [2] that the applied force at a typical point P (See Figure 2.) of a segment  $L_j$  may be written as:

$$\underline{f} = A \underline{a}_N + B |\underline{v}_N| \underline{v}_N + C |\underline{v}_T| \underline{v}_T + \underline{w} + \underline{b}_N \quad (4)$$

where  $\underline{v}_N$  is the normal component of the fluid velocity relative to the cable segment at P,  $\underline{a}_N$  is the normal component of the fluid acceleration relative to the cable segment at P,  $\underline{v}_T$  is the tangential component of the fluid velocity relative to the cable segment at P,  $\underline{w}$  is the segment weight per unit length, and  $\underline{b}_N$  is the buoyancy force per unit length.

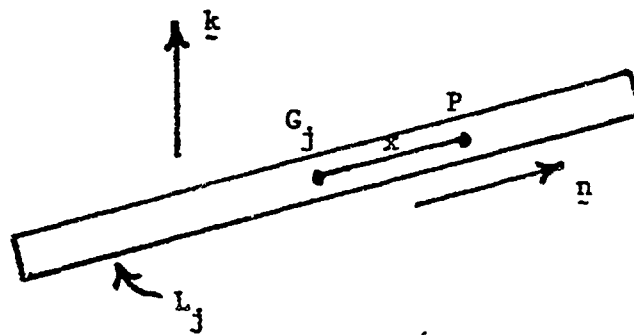


Figure 2. A Typical Cable Segment  $L_j$ .

The coefficients A, B, and C are:

$$A = C_M \rho (\pi/4) d^2 \quad (5)$$

$$B = C_N \rho (d/2) \quad (6)$$

$$C = C_T \rho (d/2) \quad (7)$$

where  $\rho$  is the fluid mass density and  $d$  is the diameter of  $L_j$ .  $C_M$ ,  $C_N$ , and  $C_T$  are coefficients dependent upon the Reynolds number of the fluid flow relative to the cable segment. They are usually determined experimentally and reported results may vary slightly. Webster [7] records them as

$$C_M = 1.0 \quad (8)$$

$$C_N = \begin{cases} 0.0 & \text{for } R_{eN} \leq 0.1 \\ 0.45 + 5.93/(R_{eN})^{0.33} & \text{for } 0.1 < R_{eN} \leq 400.00 \\ .27 & \text{for } 400 < R_{eN} \leq 10^5 \\ (0.3 & \text{for } R_{eN} > 10^5 \end{cases} \quad (9)$$

and

$$C_T = \begin{cases} 1.88/(R_{eT})^{0.74} & \text{for } 0.1 < R_{eT} \leq 100.55 \\ 0.062 & \text{for } R_{eT} > 100.55 \end{cases} \quad (10)$$

where the Reynolds numbers  $R_{eN}$  and  $R_{eT}$  are defined as:

$$R_{eN} = \rho d |\underline{v}_N| / \mu \quad (11)$$

and

$$R_{eT} = \rho d |\underline{v}_T| / \mu \quad (12)$$

where  $\mu$  is the viscosity of the fluid.

The weight force  $\underline{w}$  may be expressed as:

$$\underline{w} = -\rho_c g \underline{k} \quad (13)$$

where  $\rho_c$  is the cable mass per unit length,  $g$  is the gravity constant, and  $\underline{k}$  is a vertical (up) unit vector.

The buoyancy force at  $P$  may be expressed as [2]

$$\underline{b}_N = \rho g a \underline{n} (\underline{k} \times \underline{n}) \quad (14)$$

where  $a$  is the cross section area of  $L_j$  and  $\underline{n}$  is a unit vector parallel to  $L_j$ . (See Figure 2.) (Note that since the cable segment is a model of a portion of a continuous cable, the ends of the segment are not exposed to the fluid. That is, the segment ends are "shielded" from the fluid by the adjoining cable segments. This results in  $\underline{b}_N$  being normal to  $L_j$ .)

If the set of applied forces at all points along  $L_j$  are replaced by a single force  $\underline{F}_j$  passing through  $G_j$  together with a couple with torque  $\underline{T}_j$ , then  $\underline{F}_j$  and  $\underline{T}_j$  may be expressed as:

$$\underline{F}_j = \underline{F}_M + \underline{F}_N + \underline{F}_T + \underline{W} + \underline{B}_N \quad (15)$$

and

$$\underline{T}_j = \underline{T}_M + \underline{T}_N \quad (16)$$

where  $\underline{F}_M$ ,  $\underline{F}_N$ ,  $\underline{F}_T$ ,  $\underline{W}$ , and  $\underline{B}_N$  are due to the added mass, normal drag, tangential drag, gravity and buoyancy forces respectively, and  $\underline{T}_M$  and  $\underline{T}_N$  are torques due to the added mass and normal drag forces. These forces and torques may be expressed as [2]:

$$\underline{F}_M = AL_a \underline{G} \underline{N} \quad (17)$$

$$\begin{aligned} \underline{F}_N = & B \{ \underline{\omega} \times \underline{n} (1/3c) (X^{3/2}|_{-L/2} - X^{3/2}|_{L/2}) + [(b/2c) \underline{\omega} \times \underline{n} \\ & + \underline{v}_{GN}] \left[ \left( \frac{cL+b}{4c} \right) X^{1/2}|_{L/2} + \left( \frac{cL-b}{4c} \right) X^{1/2}|_{-L/2} \right. \\ & \left. + \left( \frac{4ac-b^2}{8c^{3/2}} \right) \log \left( \frac{X^{1/2}|_{L/2} + (Lc^{1/2}/2) + (b/2c^{1/2})}{X^{1/2}|_{-L/2} - (Lc^{1/2}/2) + (b/2c^{1/2})} \right) \right] \} \end{aligned} \quad (18)$$

$$\underline{F}_T = CL|\underline{v}_{GT}|\underline{v}_{GT} \quad (19)$$

$$\underline{W} = \underline{\omega}L = -\rho_c g L \underline{k} \quad (20)$$

$$\underline{B}_N = \underline{b}_N L = \rho g a L \underline{n} \times (\underline{k} \times \underline{n}) \quad (21)$$

$$\underline{T}_M = -A(L^3/12)[\underline{a} - (\underline{a} \cdot \underline{n})\underline{n} + (\underline{\omega} \cdot \underline{n})\underline{n} \times \underline{\omega}] \quad (22)$$

and

$$\begin{aligned} \underline{T}_N = \underline{B}_N \times \underline{v}_{GN} & \left\{ (1/3c)X^{3/2} \Big|_{L/2} - (1/3c)X^{3/2} \Big|_{-L/2} - \frac{b(cL+b)}{8c^2} X^{1/2} \Big|_{L/2} \right. \\ & + \frac{b(-cL+b)}{8c^2} X^{1/2} \Big|_{-L/2} - \frac{b(4ac-b^2)}{16c^{5/2}} \log \left( \frac{X^{1/2} \Big|_{L/2} + (Lc^{1/2}/2) + (b/2c^{1/2})}{X^{1/2} \Big|_{-L/2} - (Lc^{1/2}/2) + (b/2c^{1/2})} \right) \\ & + B[(\underline{\omega} \cdot \underline{n})\underline{n} - \underline{\omega}] \left\{ (1/4c)[(L/2) - (5b/6c)]X^{3/2} \Big|_{L/2} \right. \\ & + (1/4c)[(L/2) + (5b/6c)]X^{3/2} \Big|_{-L/2} \\ & + [5b^2 - 4ac]/16c^2 \left[ \frac{cL+b}{4c} X^{1/2} \Big|_{L/2} + \frac{cL-b}{4c} X^{1/2} \Big|_{-L/2} \right. \\ & \left. \left. + \left( \frac{4ac-b^2}{8c^{3/2}} \right) \log \left( \frac{X^{1/2} \Big|_{L/2} + (Lc^{1/2}/2) + (b/2c^{1/2})}{X^{1/2} \Big|_{-L/2} - (Lc^{1/2}/2) + (b/2c^{1/2})} \right) \right] \right\} \quad (23) \end{aligned}$$

where  $L$  is the length of  $L_j$ ,  $\underline{\omega}$  and  $\underline{a}$  are the angular velocity and angular acceleration of  $L_j$ , and  $\underline{v}_{GN}$  and  $\underline{a}_{GN}$  are the components of the fluid velocity and acceleration normal to  $L_j$  at  $G_j$ .  $\underline{v}_{GN}$  and  $\underline{a}_{GN}$  may be expressed as:

$$\underline{v}_{GN} = \underline{v}_{W/G} - (\underline{v}_{W/G} \cdot \underline{n})\underline{n} \quad (24)$$

and

$$\underline{a}_{GN} = \underline{a}_{W/G} - (\underline{a}_{W/G} \cdot \underline{n})\underline{n} \quad (25)$$

where  $\underline{v}_{W/G}$  and  $\underline{a}_{W/G}$  are the velocity and acceleration of the fluid relative to  $G_j$ . Finally, the coefficients  $a$ ,  $b$ , and  $c$ , and  $X$  are:

$$a = \underline{v}_{GN} \cdot \underline{v}_{GN} \quad (26)$$

$$b = -2\underline{v}_{GN} \cdot \underline{\omega} \times \underline{n} \quad (27)$$

$$c = (\omega \times r) \cdot (\omega \times n) \quad (28)$$

and

$$X = a + bx + cx^2 \quad (29)$$

where  $x$  is the length variable shown in Figure 2.

Using these results, the generalized active forces  $F_l$  are:

$$F_l = v_{jlk} F_{jk} + \omega_{jlk} T_{jk} \quad (30)$$

where  $F_{jk}$  and  $T_{jk}$  are the  $n_k$  components of  $F_j$  and  $T_j$  of Equations (15) and (16). (As before, there is a sum in Equation (30) over the repeated indices for the range of the indices.)

### III. SUBMERGED CABLE DYNAMICS

Analytical validation of the finite segment model (UCIN-CABLE) without fluid forces, has been obtained by comparing results predicted by the model with results obtained using other methods. Reference [3] records details of this validation. To obtain additional validation, and to develop further applications, a series of submerged cable maneuvers were simulated. These included buoy relaxations and anchor drops in water depicted in Figure 3.

In each maneuver, the cable length was 72 in. (1.83 m). Silicon rubber and nylon cables were simulated. The silicon rubber cable had a diameter of 0.163 in. (4.14 mm) and the nylon cable had a diameter of 0.1 in. (2.54 mm). The weight densities of the silicon rubber cables were 0.137 lb/in (24.0 N/m) and 0.108 lb/in (18.92 N/m) in air. (The heavier cable included a wire conductor.) The weight density of the nylon cable was 0.0448 lb/in (7.85 N/m) in air. The buoys and anchors were 2 in. (50.8 mm) diameter spheres weighing 0.025 lb (0.111 N) and 0.246 lb (1.09 N) in air respectively.

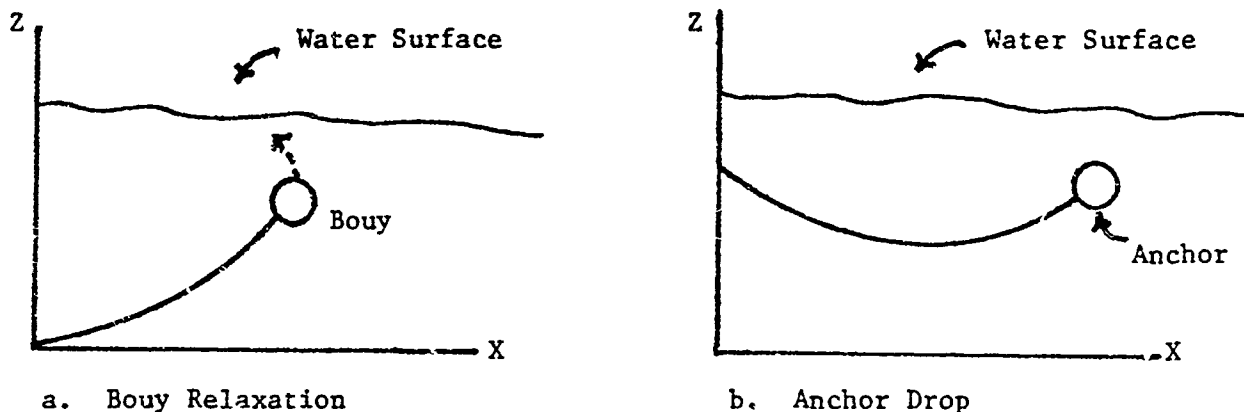


Figure 3. Cable Maneuvers.

Six tests were conducted. The first four were buoy relaxations as in Figure 3a. The last two were anchor drops as in Figure 3b. Twelve identical cable segments together with a sphere representing the buoy or anchor, were used in the model, for each test. The fluid mass density and viscosity were given the values  $1.9856 \text{ slug/ft}^3$  ( $1024.16 \text{ kg/m}^3$ ) and  $3.516 \times 10^{-5} \text{ slug/ft. sec.}$  ( $1.684 \times 10^{-3} \text{ kg/m sec}$ ) to simulate seawater.

The tests were designed to simulate experimental tests conducted at the Civil Engineering Laboratory in Port Hueneme, CA as recorded by Palo [6]. Table I provides a description of the tests. Figures 4. to 9. show the cable configuration at various times, together with comparisons with experimental results and finite element results (SEADYN) for the respective tests. Figures 10. to 15. show analogous results for the fixed end tension. Finally, Figures 16. to 21. show buoy and anchor velocities.

Test Number	Type	Initial Position of Buoy or Anchor (See Figure 3.)	Cable Material
1	Buoy Relaxation	$x = 51 \text{ in.}, z = -24 \text{ in.}$	Silicon Rubber
2	Buoy Relaxation	$x = 51 \text{ in.}, z = -33 \text{ in.}$	Silicon Rubber with Wire Core
3	Buoy Relaxation	$x = 51 \text{ in.}, z = -33 \text{ in.}$	Nylon
4	Buoy Relaxation	$x = 51 \text{ in.}, z = 16.6 \text{ in.}$	Silicon Rubber with Wire Core
5	Anchor Drop	$x = 66 \text{ in.}, z = 0$	Silicon Rubber
6	Anchor Drop	$x = 54 \text{ in.}, z = 0$	Silicon Rubber

Table I. Test Descriptions.

TEST 1

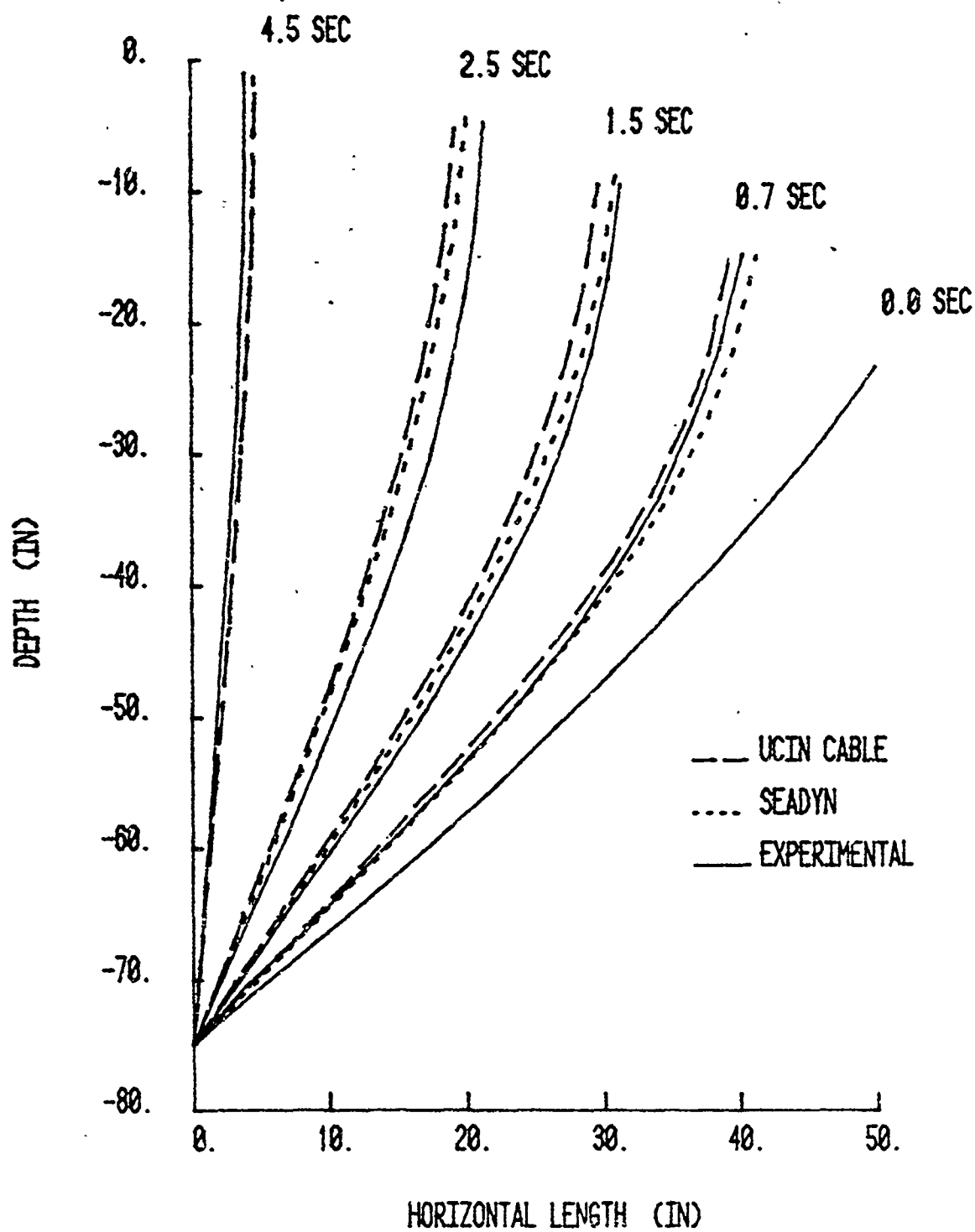


Figure 4. Test 1: Configurations for Buoy Relaxation For Rubber Cable.



# TEST 2

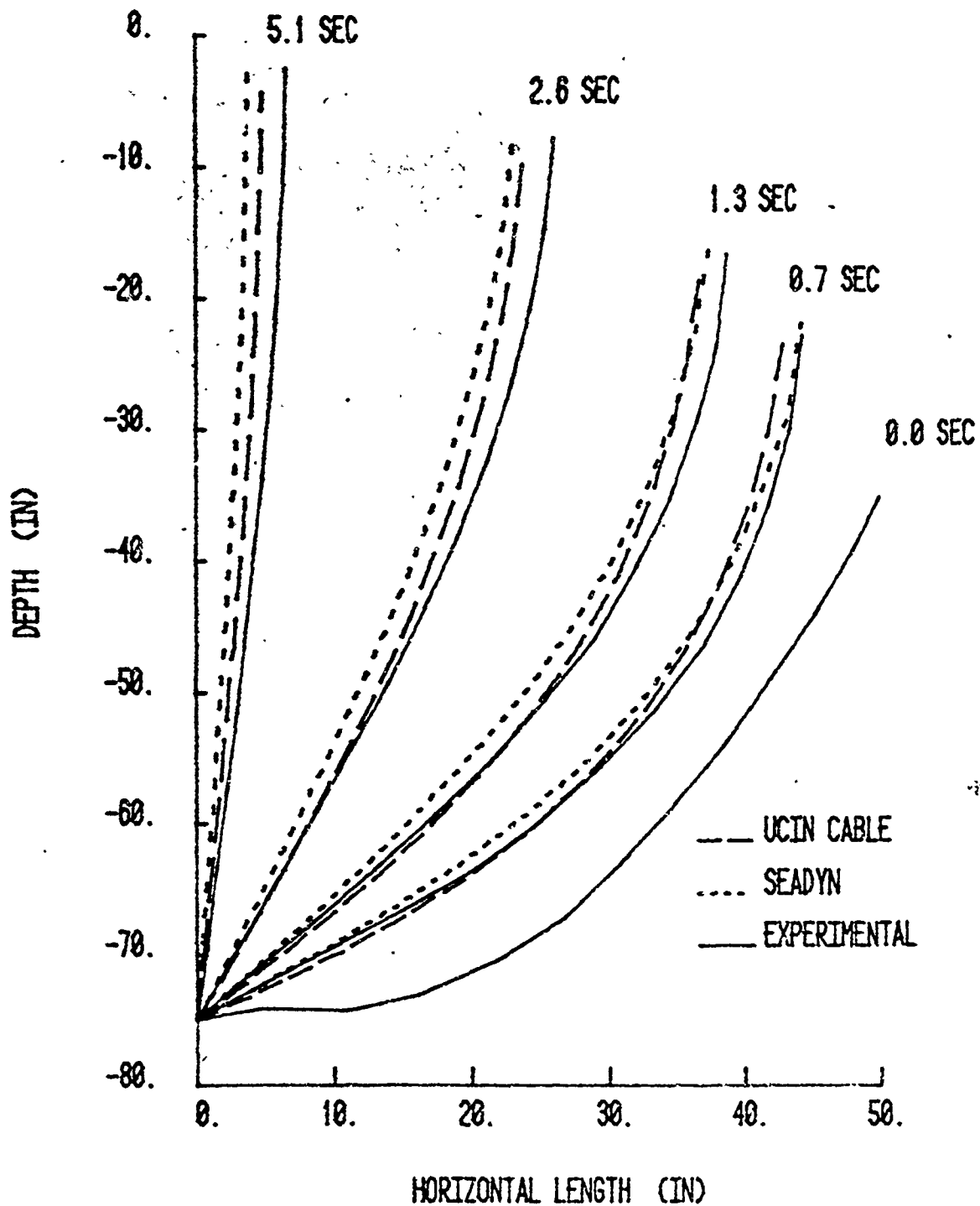


Figure 5. Test 2: Configurations for Buoy Relaxation for Rubber Cable with Wire Core.

# TEST 3

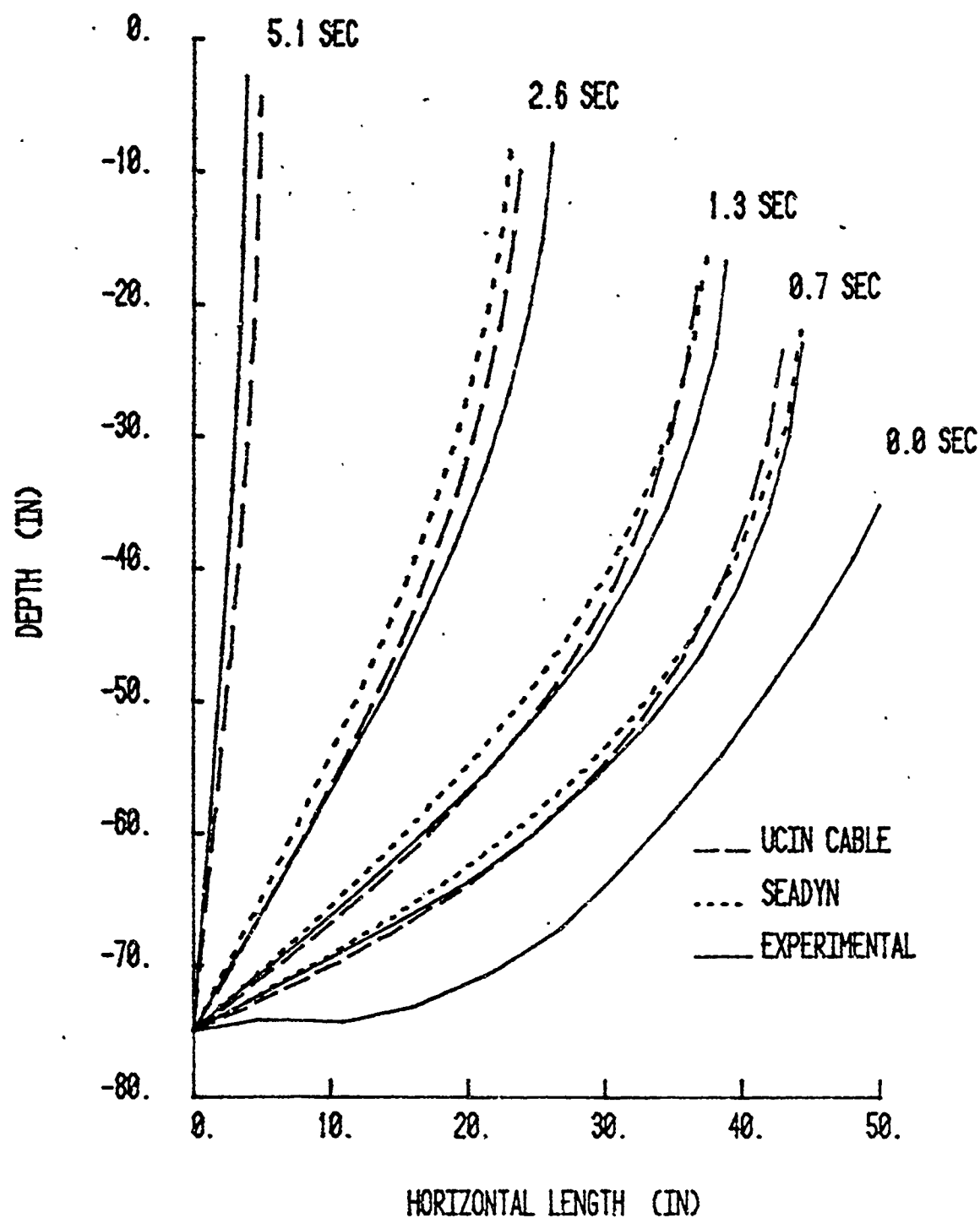


Figure 6. Test 3: Configurations for Buoy Relaxation for Nylon Cable.

# TEST 4

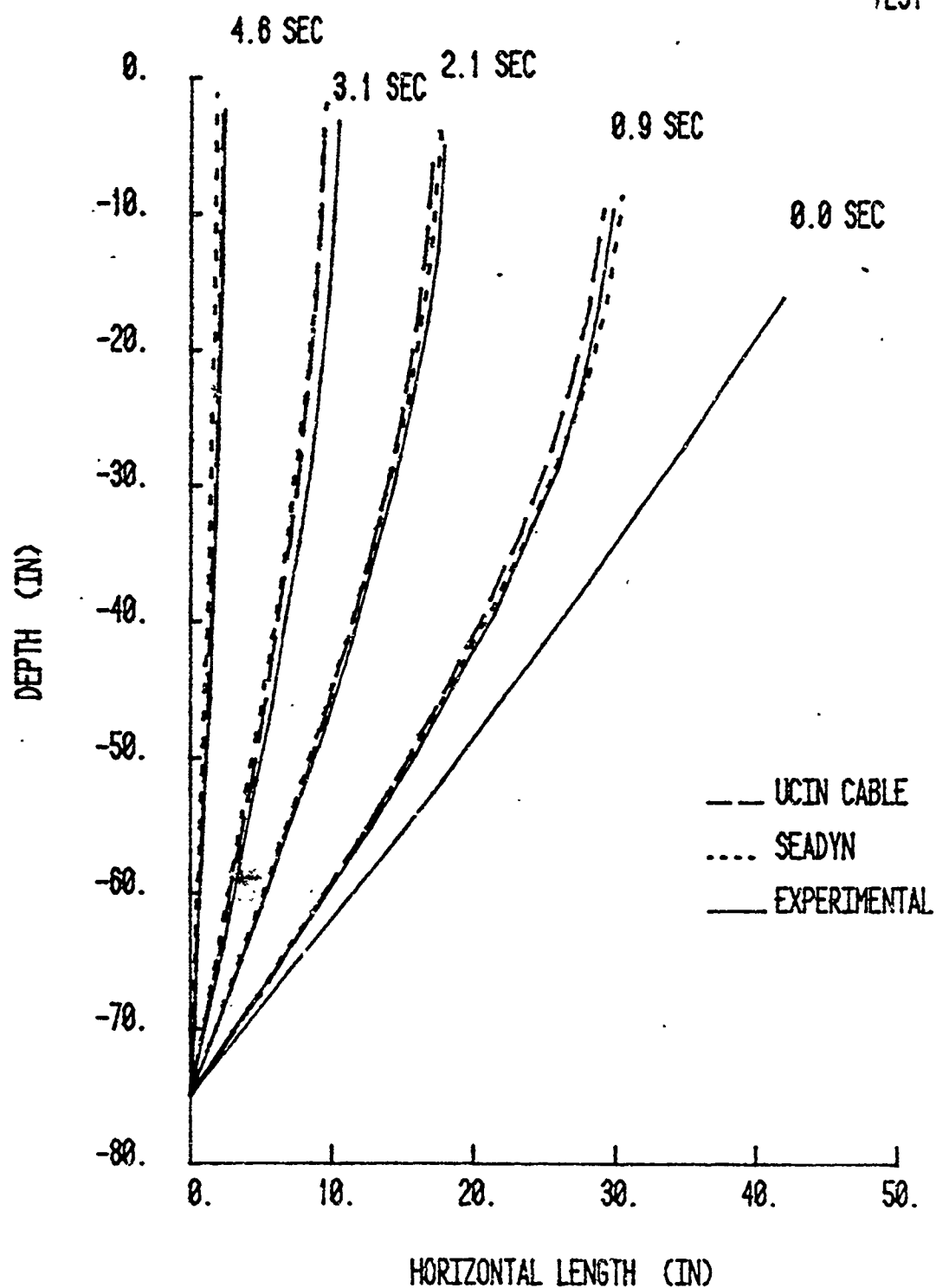


Figure 7. Test 4: Configurations for Buoy Relaxations for Rubber Cable with Wire Core.

# TEST 5

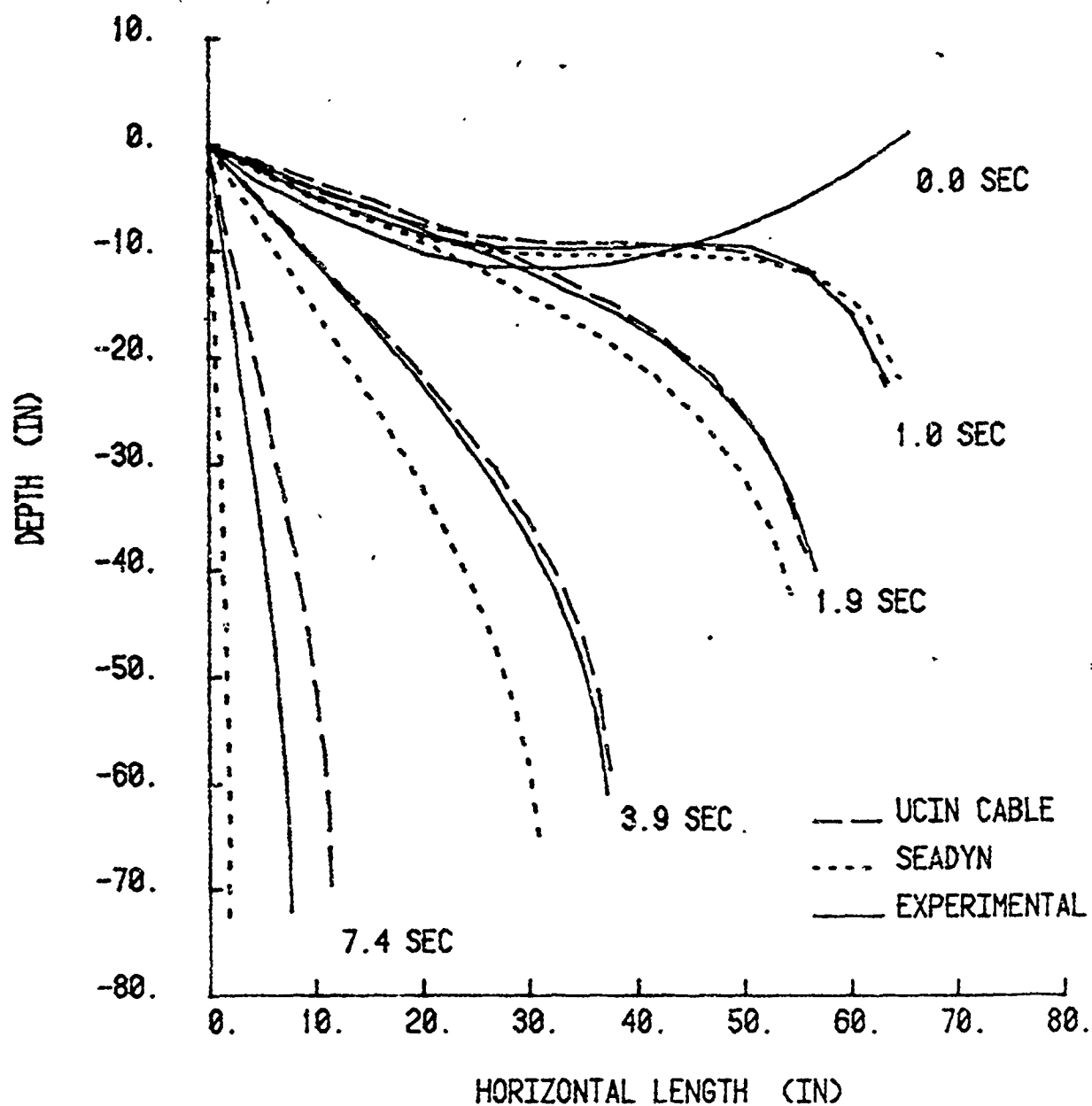


Figure 8. Test 5: Configurations for Anchor Drop for Rubber Cable.

# TEST 6

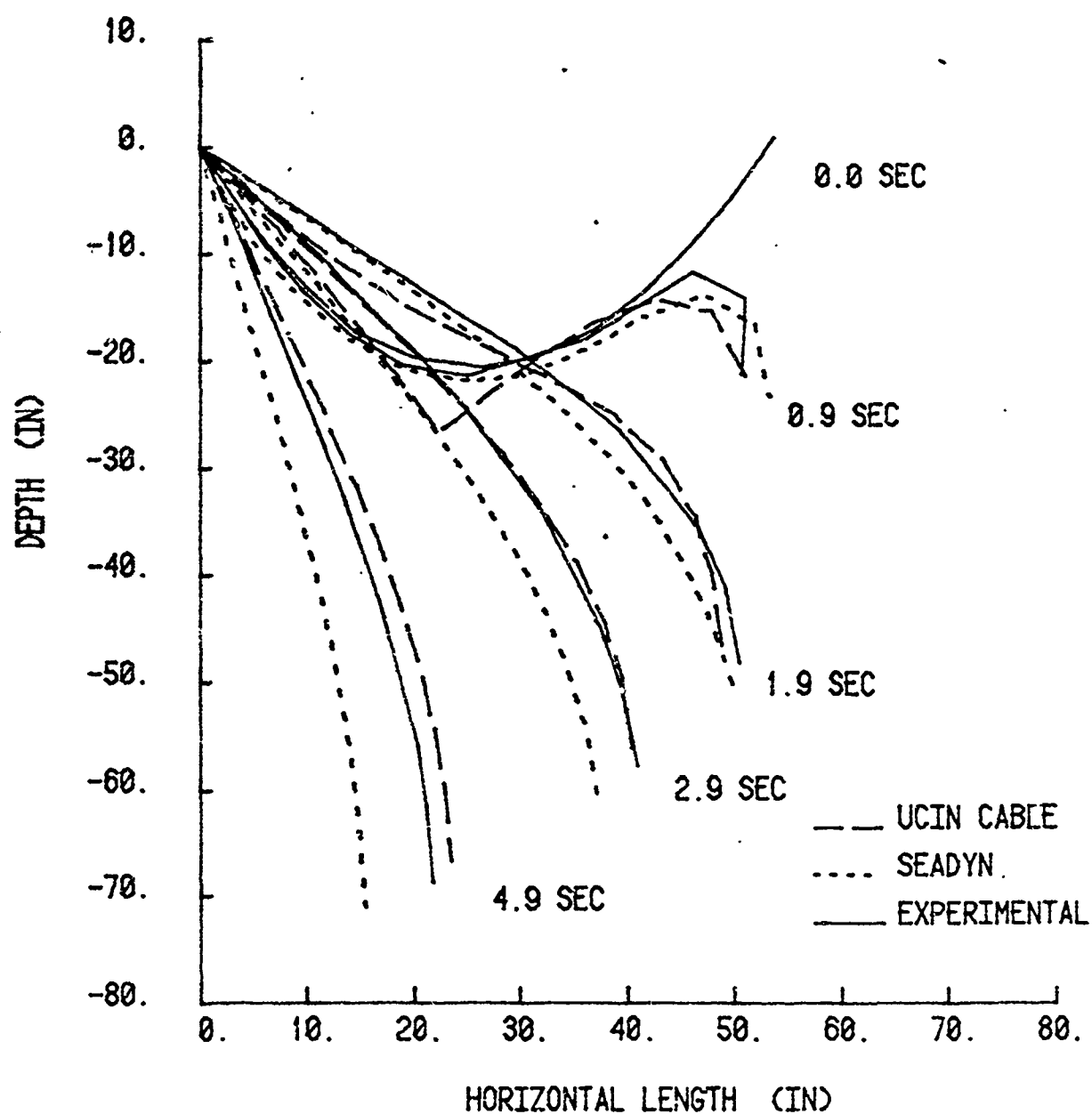


Figure 9. Test 6: Configurations for Anchor Drop for Rubber Cable.

TEST 1

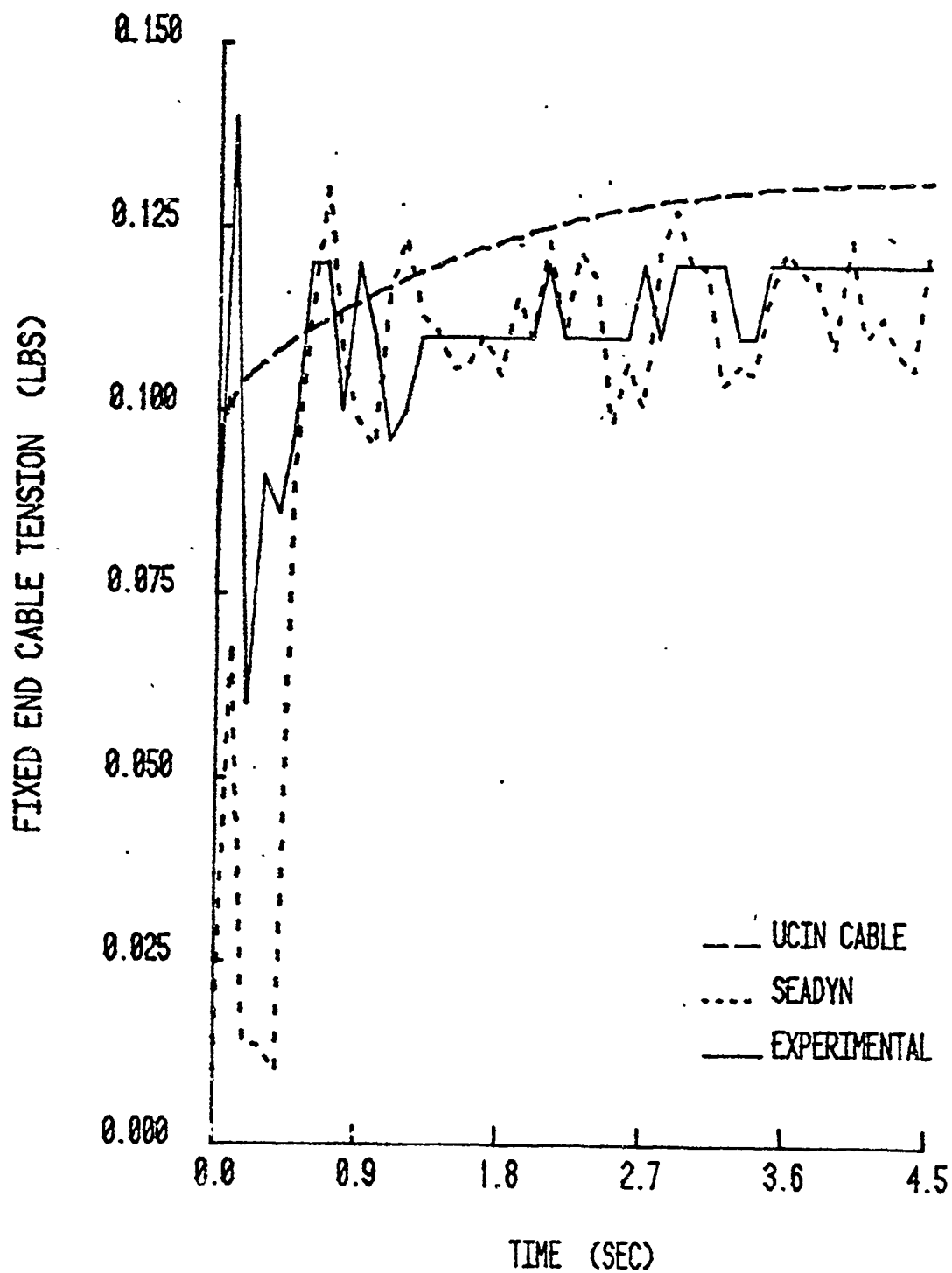


Figure 10. Test 1: Fixed End Tension for Buoy Relaxation for Rubber Cable.

# TEST 2

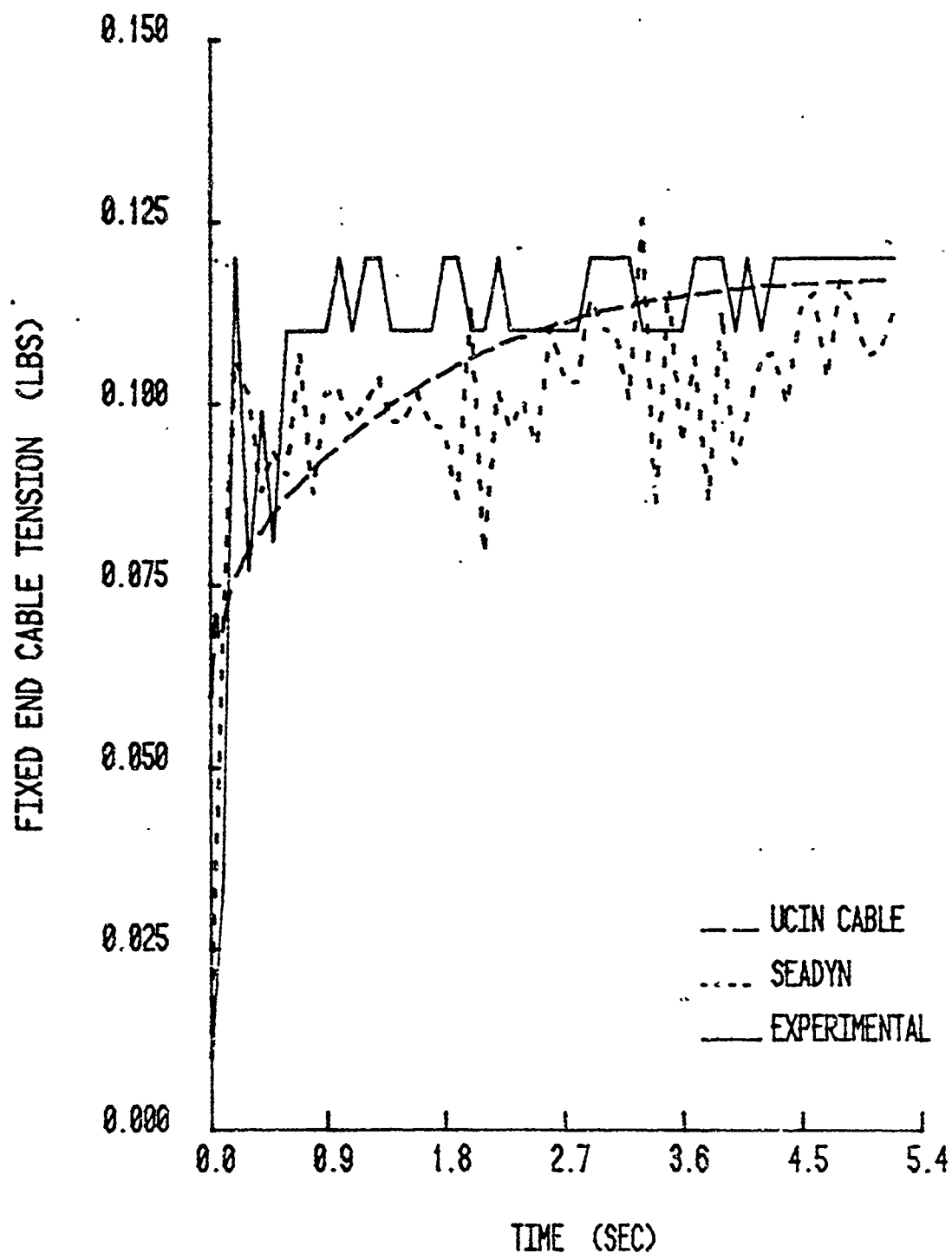


Figure 11. Test 2: Fixed End Tension for Buoy Relaxation for Rubber Cable with Wire Core.

TEST 3

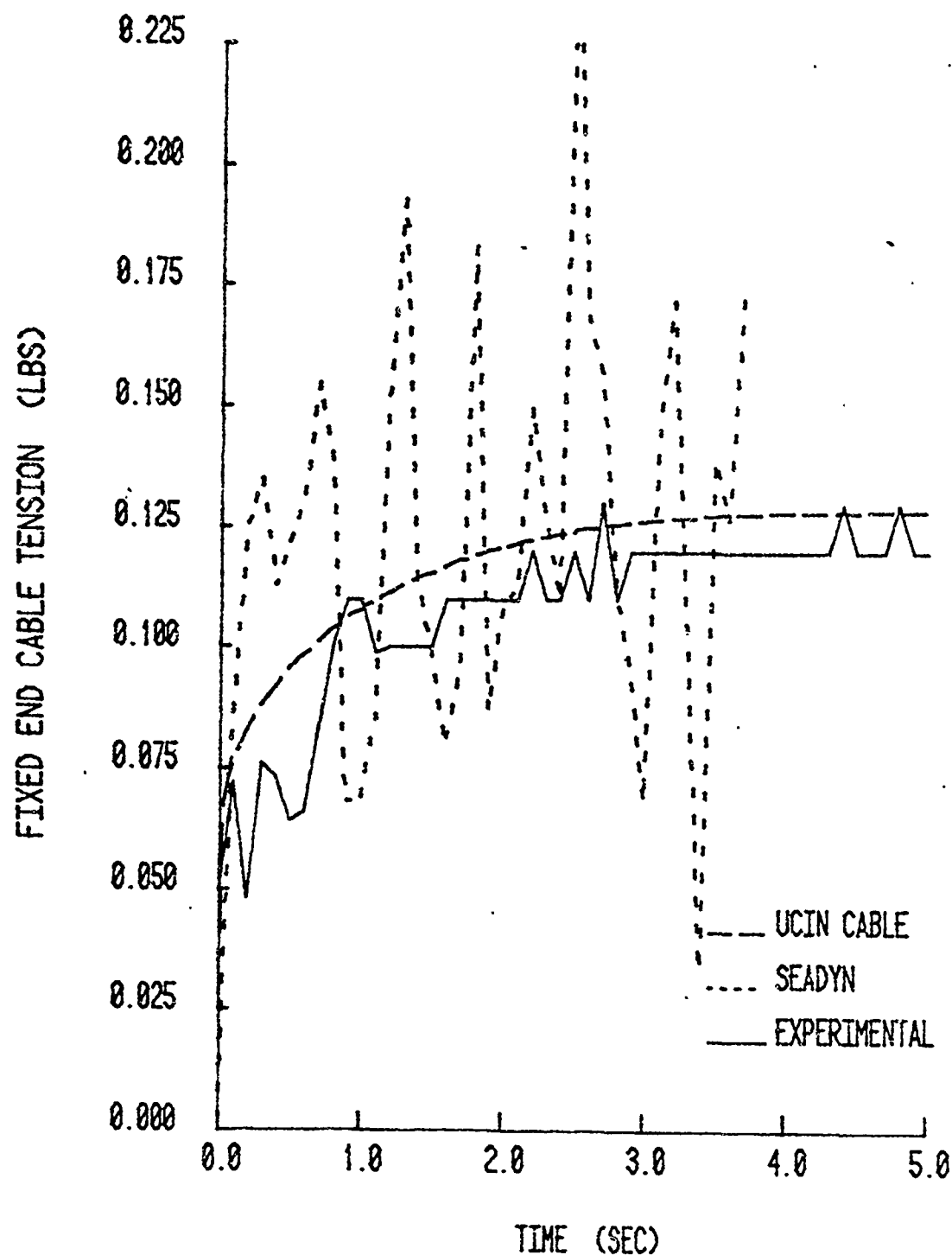


Figure 12. Test 3: Fixed End Tension for Buoy Relaxation for Nylon Cable.



TEST 4

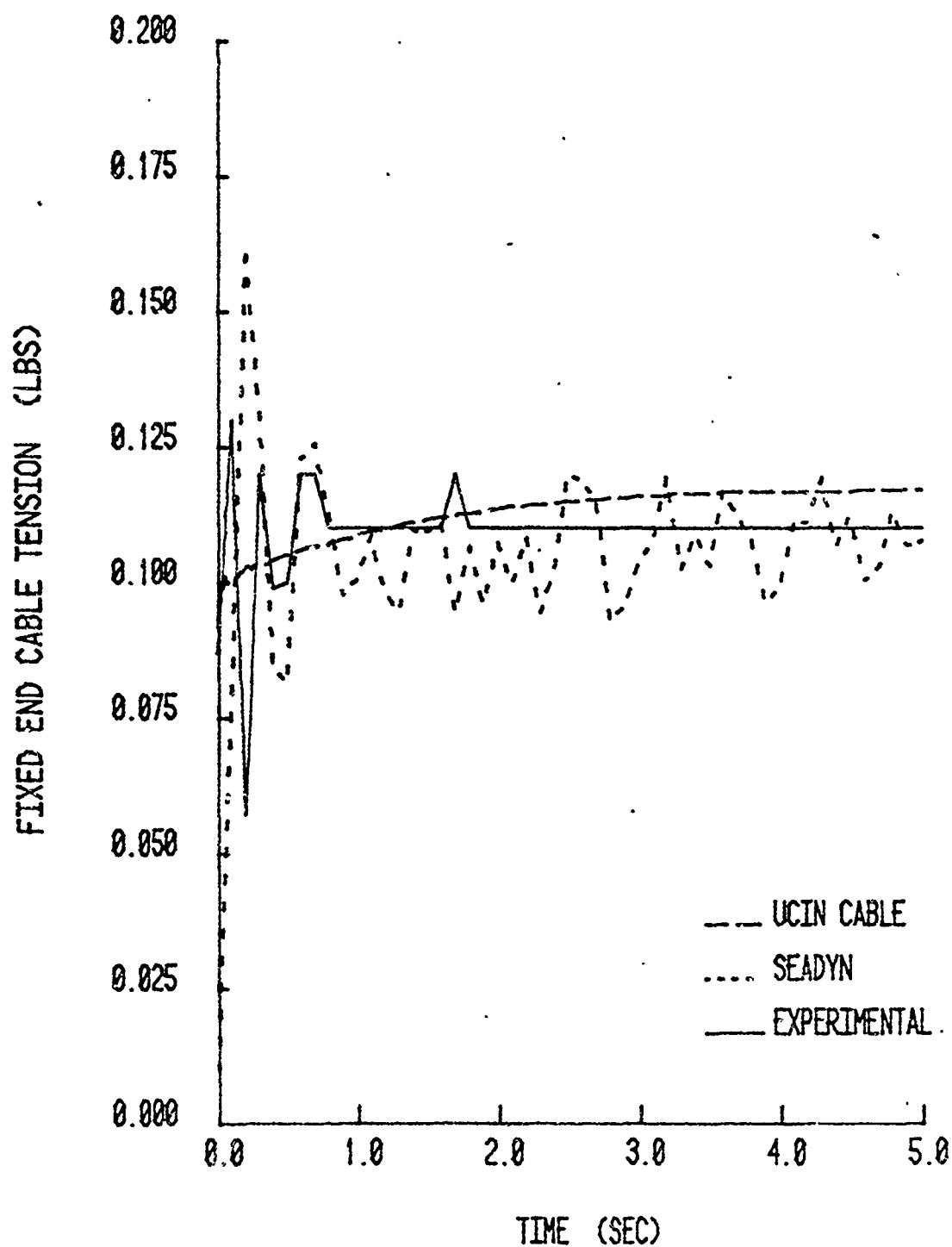


Figure 13. Test 4: Fixed End Tension for Buoy Relaxation for Rubber Cable with Wire Core.

# TEST 5

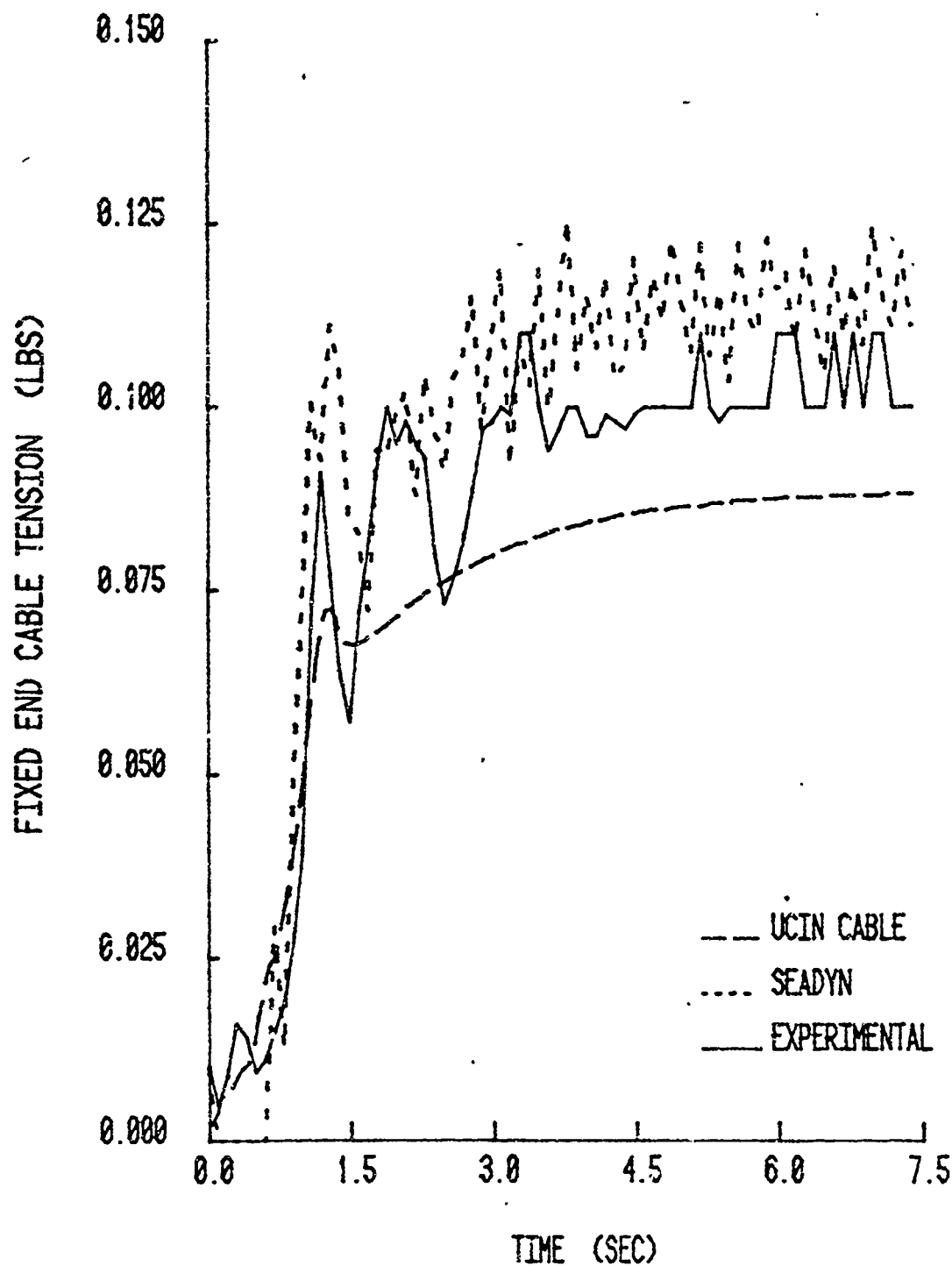


Figure 14. Test 5: Fixed End Tension for Anchor Drop for Rubber Cable.

TEST 6

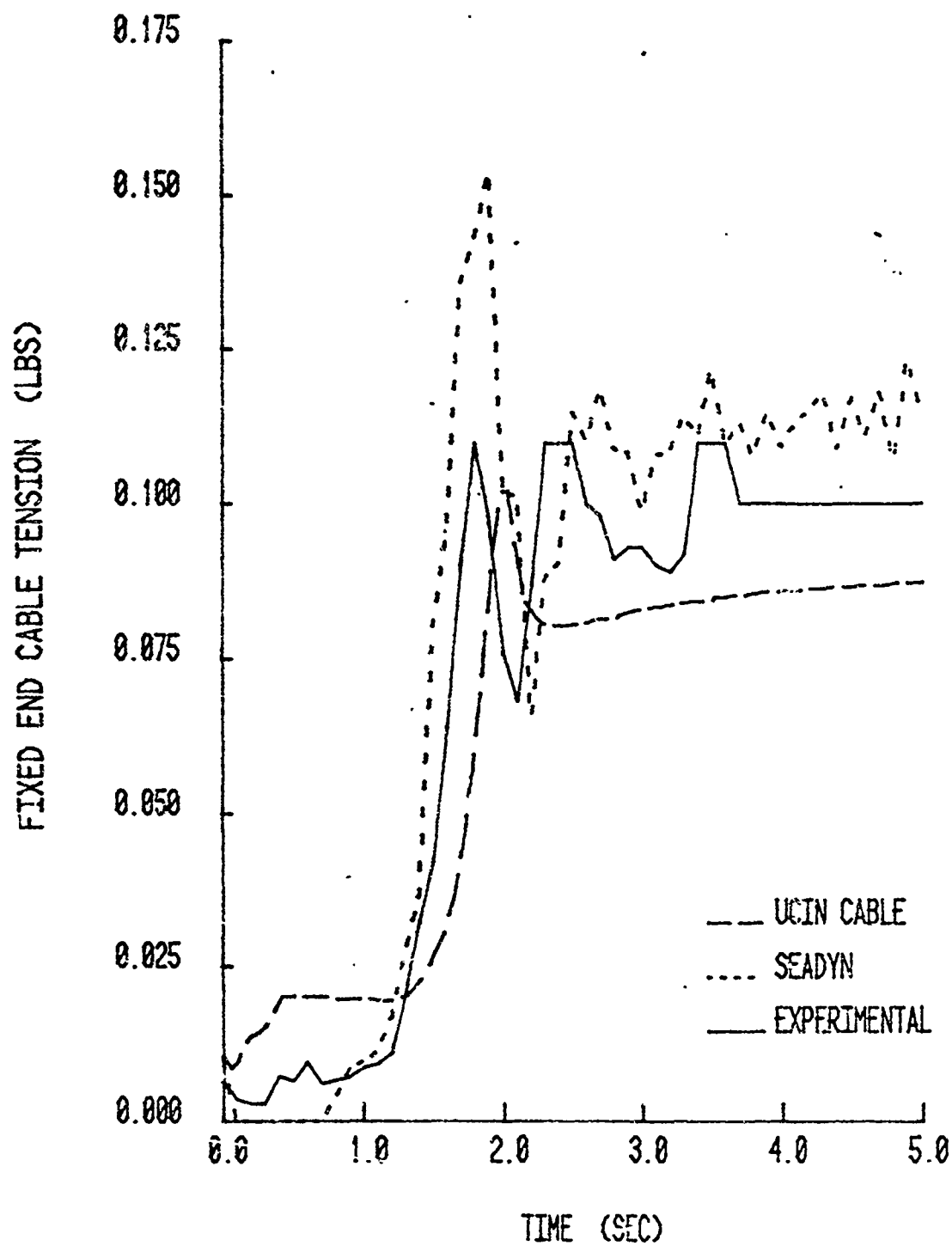


Figure 15. Test 6: Fixed End Tension for Anchor Drop for Rubber Cable.

TEST 1

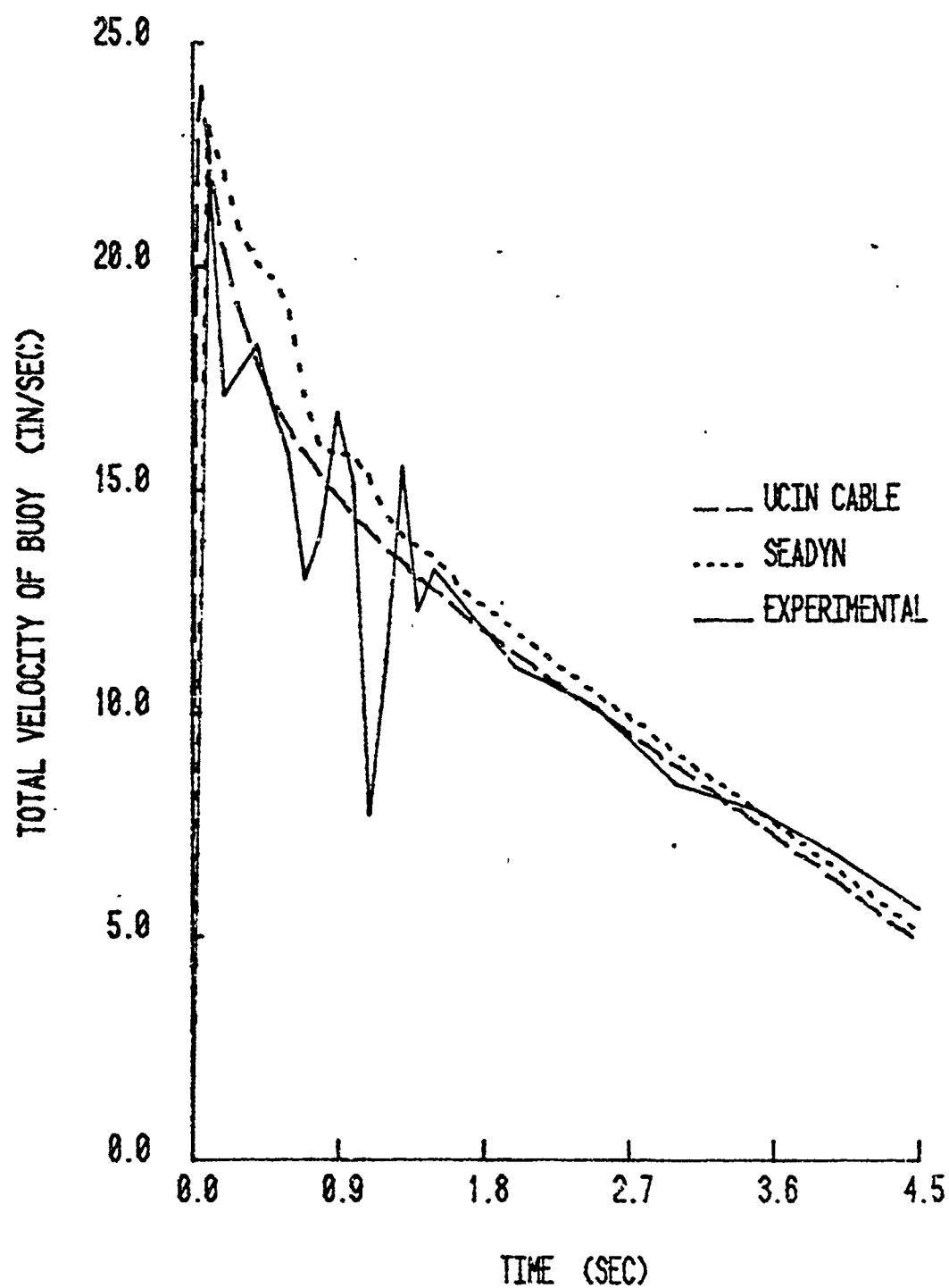


Figure 16. Test 1: Buoy Velocity for Rubber Cable.

TEST 2

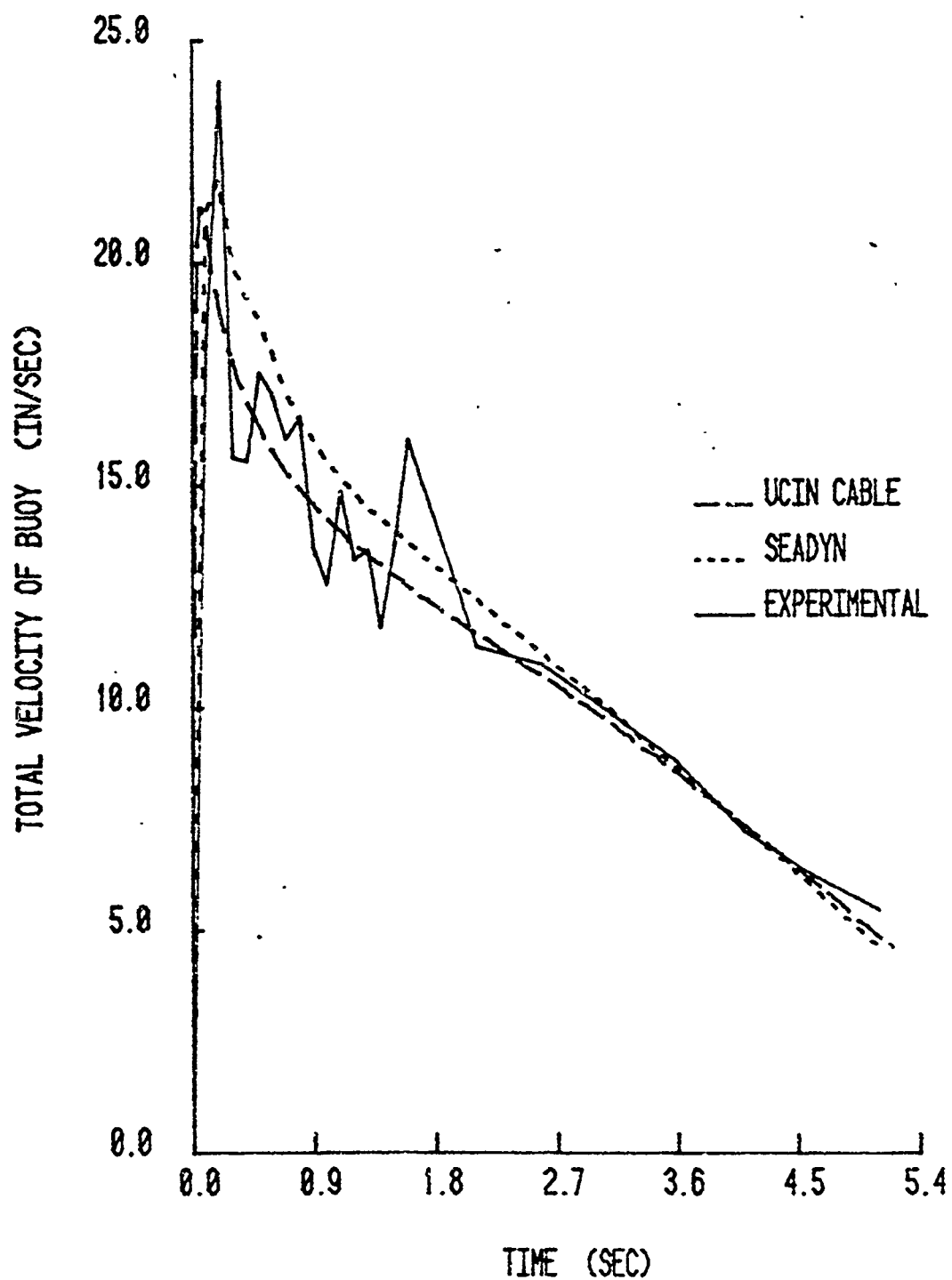


Figure 17. Test 2: Buoy Velocity for Rubber Cable with Wire Core.

TEST 3

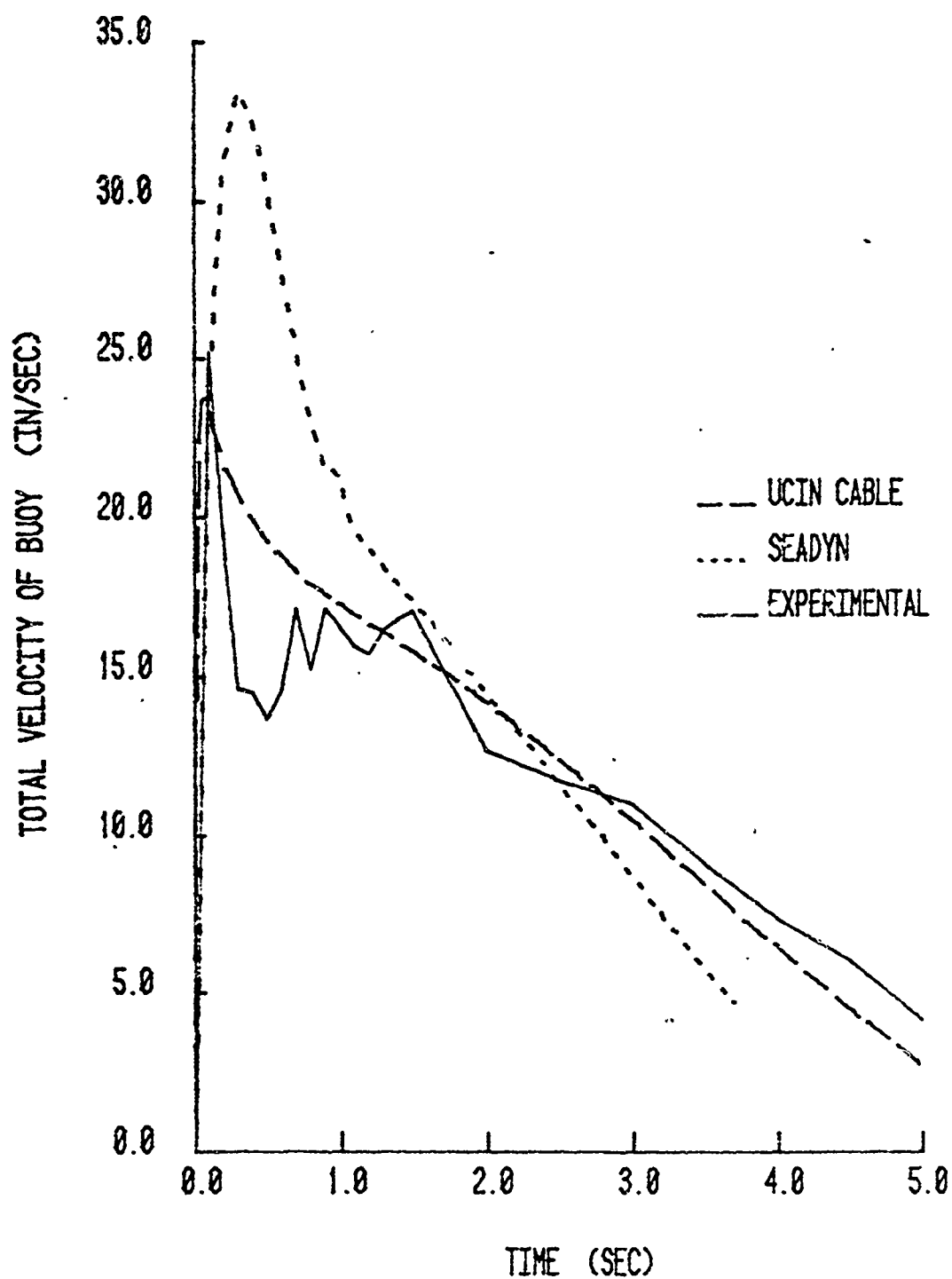


Figure 18. Test 3: Buoy Velocity for Nylon Cable.

# TEST 4

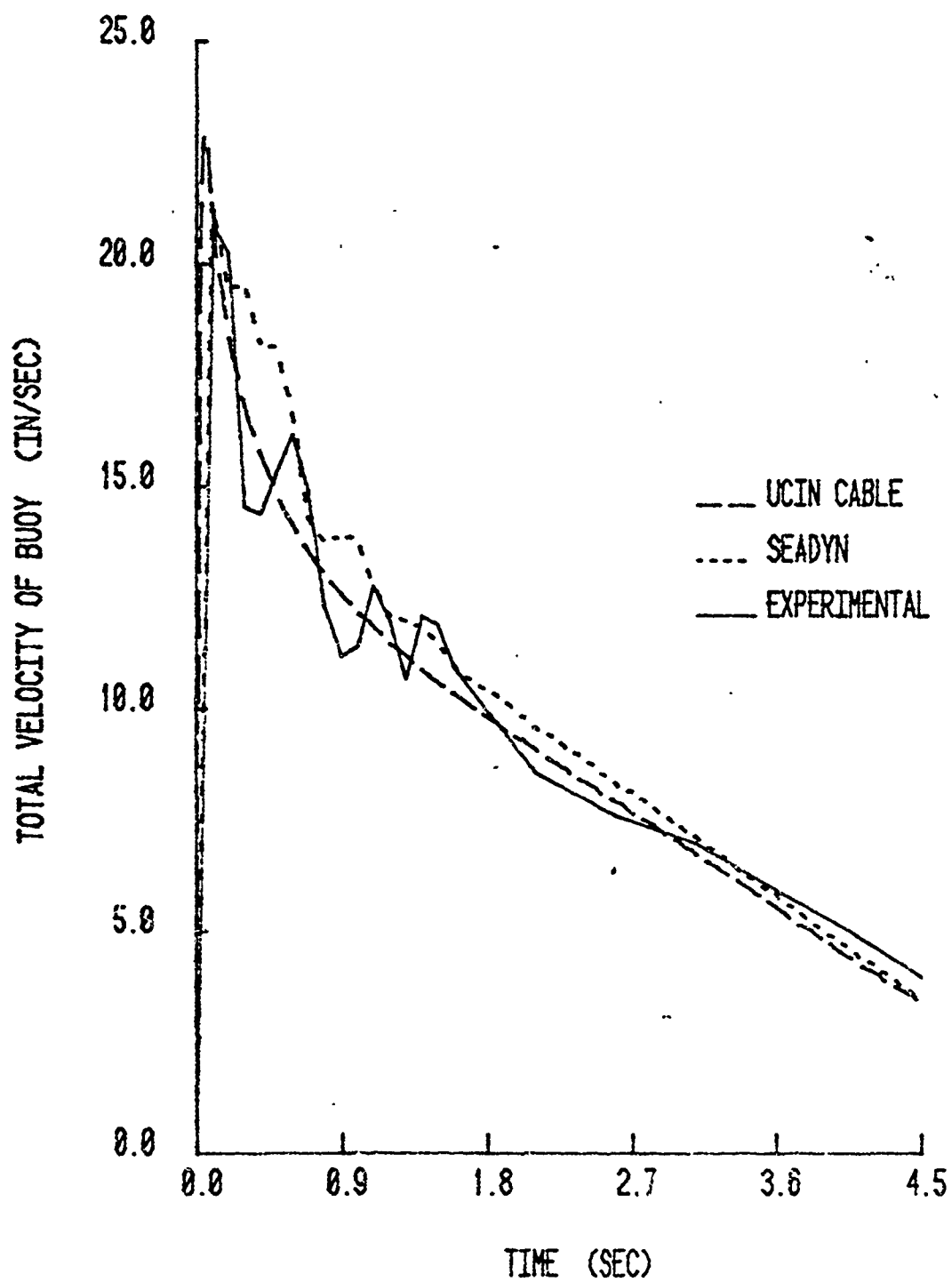


Figure 19. Test 4: Buoy Velocity for Rubber Cable with Wire Core.

# TEST 5

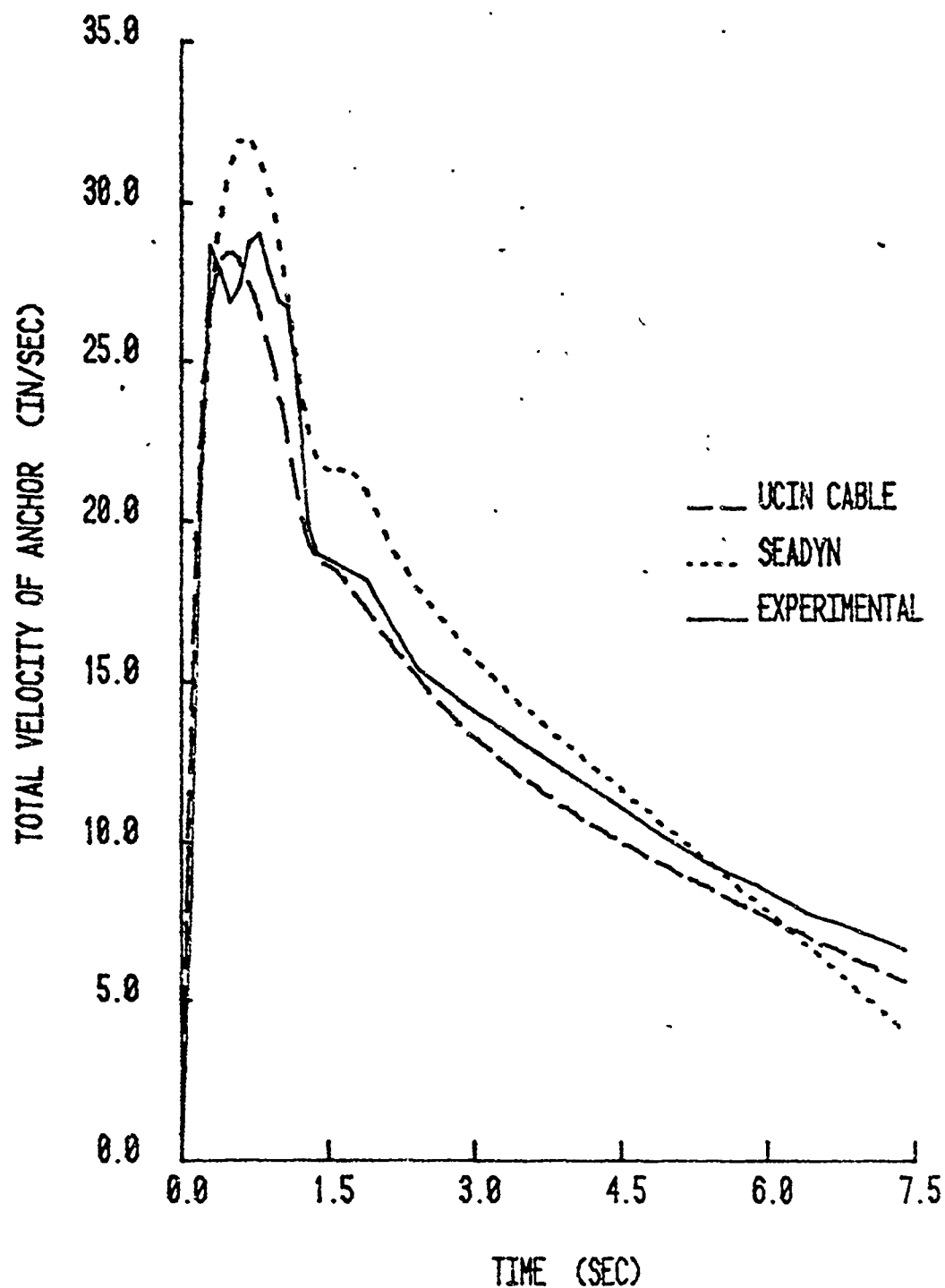


Figure 20. Test 5: Anchor Velocity for Rubber Cable.



# TEST 6

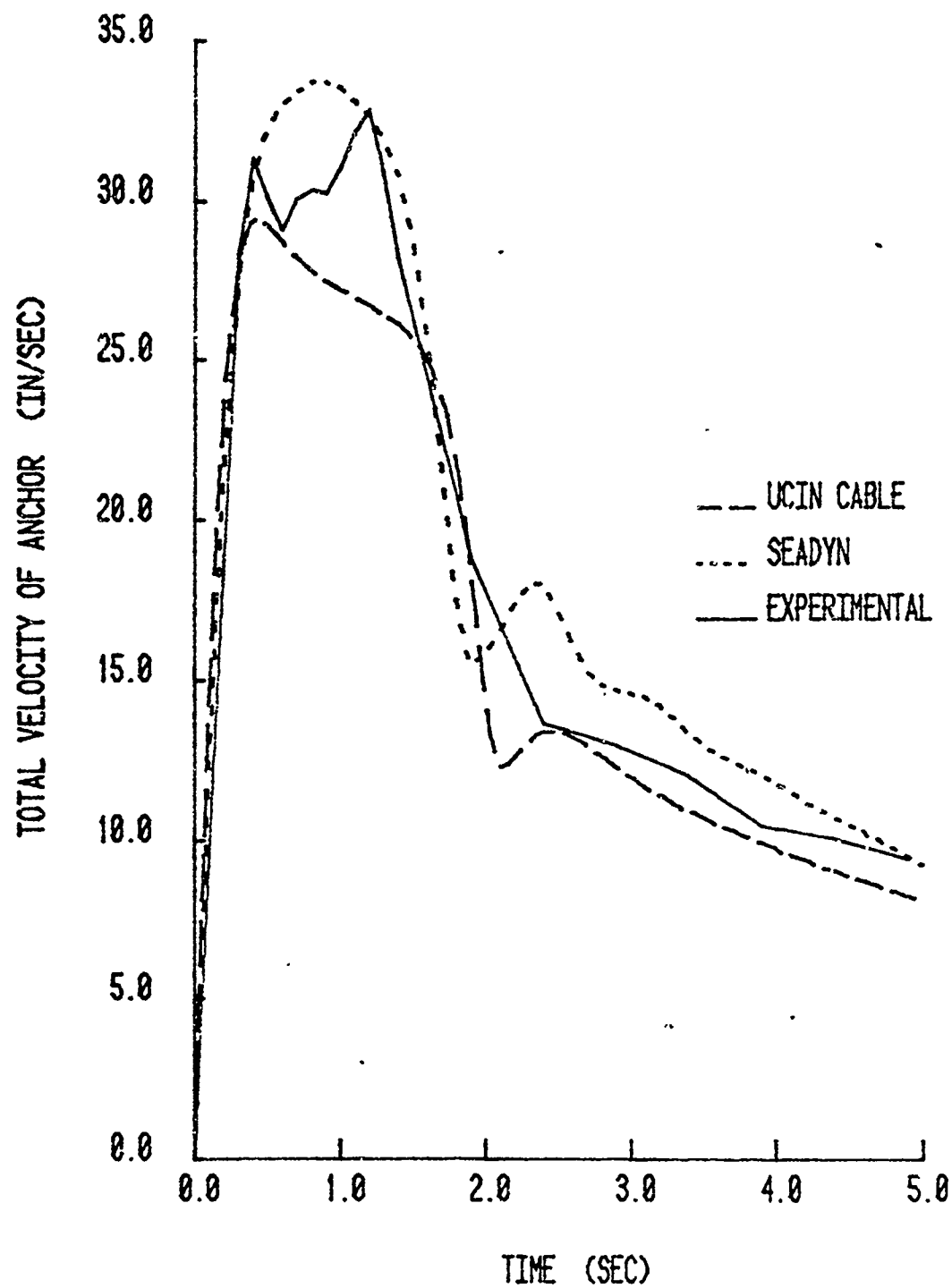


Figure 21. Test 6: Anchor Velocity for Rubber Cable.

#### IV. DISCUSSION AND CONCLUSIONS

In the current set of tests the cable was relatively light. Hence, the viscous forces were relatively large compared with the gravitational and inertia forces. The tests thus provide a validation of the fluid force modelling of the UCIN-CABLE code. (As noted earlier, the inertia force validation is reported in Reference [3].)

The results show that it is not only possible, but it is also practical to obtain numerical simulation of cable dynamics through finite segment modelling. What remains is a validation of the modelling for three-dimensional maneuvers.

# ACKNOWLEDGEMENT

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